Efficient approximation of the exponential operator for discrete 2D advection-diffusion problems

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SUMMARY

In this paper we compare Krylov subspace methods with Faber series expansion for approximating the matrix exponential operator on large, sparse, nonsymmetric matrices. We consider in particular the case of Chebyshev series, corresponding to an initial estimate of the spectrum of the matrix by a suitable ellipse. Experimental results upon matrices with large size, arising from space discretization of 2D advection-diffusion problems, demonstrate that the Chebyshev method can be an effective alternative to Krylov techniques. Copyright © 2003 John Wiley & Sons, Ltd.

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1. Introduction

Spatial discretization of linear advection-diffusion equations like

\[
\begin{align*}
  \frac{\partial u}{\partial t} & = \Delta u - \theta_1 \frac{\partial u}{\partial x} - \theta_2 \frac{\partial u}{\partial y}, \quad (x, y) \in \Omega, \; t > 0 \\
  u(x, y, 0) & = u_0(x, y), \quad (x, y) \in \Omega \\
  u(x, y, t) & \equiv 0, \quad (x, y) \in \partial \Omega, \; t > 0
\end{align*}
\]

(1)

where $\Omega \in \mathbb{R}^2$, yields naturally large systems of ordinary differential equations of the form

\[
\begin{align*}
  \dot{y}(t) & = By(t), \; t > 0 \\
  y(0) & = y_0
\end{align*}
\]

where $B$ is a large sparse nonsymmetric $n \times n$ matrix. In several papers [5, 8, 9, 11, 12, 13] explicit time integration schemes have been proposed for the solution of (2), which rest on the

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possibility of efficiently approximating the exponential operator \( \exp(A)v \), where \( v \in \mathbb{R}^n \) is a given vector, \( A = \tau B \) and \( \tau > 0 \) is a scaling factor related to the step-size \( \Delta t \).

In this framework, the Krylov subspace approach is the most popular for computing \( \exp(\tau B)v \); see, e. g., [4, 12, 14, 23, 25]. As known, the general idea of Krylov-based methods in computing \( f(A)v \) for (analytic) matrix functions, consists of approximately projecting \( f(A) \) onto a “small” Krylov subspace \( K_m = \text{Span}\{v, Av, \ldots, A^{m-1}v\} \). This reduces the problem to compute \( f(H_m) \) in view of the approximation

\[
f(A)v \approx ||v||_2 V_m f(H_m)e_1,
\]

\( V_m = [v_1, v_2, \ldots, v_m], \{v_i\} \) being the orthonormal basis of \( K_m \) and \( H_m \) the Hessenberg matrix produced by the Arnoldi process with \( v/||v||_2 \) as initial vector. When \( A \) is symmetric, the Krylov process reduces to the Lanczos method which is based on a three-term recurrence. In [1], we analyzed the Chebyshev series approximation of \( f(A) = \exp(A) \) for \( A \) symmetric negative-definite, which is on turn based on the three-term recurrence for Chebyshev polynomials on the spectral interval of \( A \) (and whose implementation does not require any scalar product). Despite of the preprocessing cost due to extremal eigenvalue approximation and of the generally higher convergence rate of the Lanczos method, the Chebyshev series approach proves an effective alternative to Krylov techniques, especially when memory bounds do not allow the storage of the whole subspace \( K_m \).

In the nonsymmetric case the Ritz vectors \( \{v_1, \ldots, v_m\} \) satisfy a long-term recurrence which leads to a \( O(m^2) \) complexity. Actual implementations of the exponential operator [23, 25] try to keep \( m \) reasonably small in order to control the quadratic cost, and possibly the storage occupancy when the problem is extremely large. This a priori choice of \( m \) entails that the time step \( \Delta t \) is dynamically subdivided into a sequence of small substeps \( \tau_i, i = 1, \ldots p \), where

\[
\exp(\Delta t B)v = \exp(\tau_p B)(\exp(\tau_{p-1} B)(\ldots (\exp(\tau_1 B)v) \ldots)),
\]

in order to attain a prescribed accuracy.

The natural extension of Chebyshev series to function of nonsymmetric matrices is given by Faber series. Given a compact set \( \Omega \subset \mathbb{C} \) including the spectrum \( \sigma(A) \) of \( A \), and assuming \( f \) analytic on a neighborhood of \( \Omega \), we can approximate \( f(A)v \) with

\[
f(A)v \approx s_m := S_m(A)v := \sum_{j=0}^{m} a_j(f)F_j(A)v,
\]

where \( F_j \) and the \( a_j \) are respectively the \( j \)th Faber polynomial and the \( j \)th Faber coefficient of the scalar function \( f(z) \), \( z \in \Omega \). It is important to notice that, when \( f \) is an entire function (as the exponential), convergence of (4) is superlinear. As in the Krylov approach, a long-term recurrence is now involved. A possibility to overcome this drawback, once the spectrum has been estimated by the convex hull of some marginal eigenvalues, consists in truncating the related Faber recurrence at a fixed length [19, 21]. Alternatively, we propose here to estimate the spectrum of \( A \) by an ellipse in which case, as it is well known, the recurrence becomes three-term and Faber polynomials become scaled and translated Chebyshev polynomials. This choice is particularly well-justified in the present framework, since spectra of “advectio-diffusion” matrices, arising from stable FD discretizations, turn out to be very similar to ellipses (possibly degenerating); see the figures in §4. Our estimating ellipse is obtained cheaply via a rough
approximation of the extremal eigenvalues, using few Arnoldi iterations (within the ARPACK routines [15]).

Starting from the pioneering work by Manteuffel [16, 17] in the ’70s, it is worth recalling here the use of Chebyshev and Faber polynomials in the construction of efficient semi-iterative methods for the solution of nonsymmetric linear systems. Also in these cases, a cheap estimate of the spectrum by Arnoldi iterations represents a key ingredient; see, e. g. [6, 28]. Our numerical results show that, despite of the preprocessing stage of extremal eigenvalues evaluation, the Chebyshev approximation of the exponential matrix method is almost always faster than Krylov, on advection-diffusion discretization matrices. This is due to the need of step fractioning in Krylov implementations, and happens even when the Krylov threshold dimension \( m \) is optimally chosen, with speedup ratios up to about 3. On the other hand, the Chebyshev preprocessing stage can become negligible in practice when time discretization of (2) is considered and several matrix exponential have to be computed, since the extremal eigenvalues of \( B \) can be computed once and for all. Moreover, this stage could benefit of the growing advances in the computation of eigenvalues of nonsymmetric matrices (we quote e. g. the recent Jacobi-Davidson method [26]).

The paper is organized as follows. In \( \S2 \) we give an outline about the fundamental properties of Faber and Chebyshev polynomials and series in the complex domain, and in \( \S3 \) we present some results on error bounds for the approximation of the exponential operator. In \( \S4 \) we present and discuss our implementation of the Chebyshev algorithm. Finally, in \( \S5 \) we compare Krylov and Chebyshev methods on a set of numerical tests, concerning computation of \( \exp(\Delta t B)v \), where \( B \) comes from Finite Difference (FD) discretization of the advection-diffusion equation (1). These results show that, even more clearly than in the symmetric case [1], the Chebyshev approach is an effective alternative to Krylov methods in computing the exponential operator.

2. Background on Faber series

In this section we recall some basic classical results, concerning approximation of analytic functions on compact sets of the complex plane based on Faber series. For a throughout discussion of this topic we refer the reader e. g. to [3, 18, 27].

Let \( B[0,r] := \{ w \in \mathbb{C} : |w| \leq r \} \) and \( \overline{\mathbb{C}} := \mathbb{C} \cup \{ \infty \} \); if \( K \subset \mathbb{C} \) is a compact set with more than one point, then there is a function \( w = \phi(z) \) which maps \( \overline{\mathbb{C}} \setminus K \) conformally onto \( \overline{\mathbb{C}} \setminus B[0,1] \) and satisfies the conditions

\[
\phi(\infty) = \infty, \quad \lim_{z \to \infty} \frac{\phi(z)}{z} = \frac{1}{\gamma},
\]

where \( \gamma \) is called capacity of \( K \); moreover, given any integer \( j > 0 \), the function \( [\phi(z)]^j \) has a Laurent expansion at infinity of the form

\[
[\phi(z)]^j = \frac{1}{\gamma^j} \left( z^j + d_0^{(j)} z^{j-1} + \ldots + d_{j-1}^{(j)} + \frac{d_j}{z} + \ldots \right),
\]

(cf. [18, Theorem 3.14 and Corollary]). The polynomial

\[
F_j(z) := \frac{1}{\gamma^j} (z^j + d_0^{(j)} z^{j-1} + \ldots + d_{j-1}^{(j)})
\]
is called \( j \)th \textit{Faber polynomial} generated by \( K \). Now, let \( z = \psi(w) \) the inverse function of \( \phi(z) \); its Laurent expansion at infinity takes the form

\[
\psi(w) = \gamma \left( w + c_0 + \frac{c_1}{w} + \ldots \right) .
\]  

(6)

Then, if we set \( F_0(z) := 1 \), we get the following recurrence for the Faber polynomials (cf. [3])

\[
F_0(z) = 1, \quad F_1(z) = \frac{z}{\gamma} - c_0, \quad F_j(z) = F_1(z) \cdot F_{j-1}(z) - (c_1 F_{j-2}(z) + c_2 F_{j-3}(z) + \ldots + c_{j-1} F_0(z)) - (j - 1)c_{j-1}, \quad j \geq 2.
\]

(7)

For any \( R \geq \gamma \), let \( \Gamma_R := \{ z : |\phi(z)| = R/\gamma \} \) and \( K_R \) the bounded domain with boundary \( \Gamma_R \) (\( K_R \supseteq K \gamma \equiv K \)): from [18, Theorem 3.17] we know that every analytic function \( f \) on \( \text{int}_C K_{R_1} \) \( (R_1 > \gamma) \) can be expanded into the \textit{Faber series}

\[
f(z) = \sum_{j=0}^{\infty} a_j(f) F_j(z),
\]

(8)

where

\[
a_j(f) := \frac{1}{2\pi i} \int_{|w|=R_0} \frac{f(\psi(w))}{w^{j+1}} \, dw, \quad (\gamma < R_0 < R_1)
\]

(9)

are the \textit{Faber coefficients}.

When \( f(\cdot) \equiv \exp(\cdot) \), due to its analyticity on the whole complex plane, the right-hand side of (8) converges uniformly on every compact subset of \( C \) (cf. [27, §2.1.3]). Moreover, setting \( S_m(z) := \sum_{j=0}^{m} a_j(\exp) F_j(z) \), we know from [7] that \( \{S_m\}_m \) is \textit{maximally convergent} to the exponential function in \( K \), that is

\[
\lim_{m \to \infty} \| \exp(\cdot) - S_m(\cdot) \|_{K}^{1/m} = \lim_{m \to \infty} \| \exp(\cdot) - p^*_m(\cdot) \|_{K}^{1/m},
\]

where \( p^*_m(\cdot) \) is the best uniform approximation polynomial of degree \( m \) for \( \exp(\cdot) \) on \( K \) and \( \| \cdot \|_{K} \) denotes the maximum norm on \( K \).

From these properties, a polynomial approximation method for the \textit{matrix exponential operator} can be derived. Indeed, given \( K \) compact set, \( K \supseteq \sigma(A) \), considering the matrix polynomial

\[
S_m(A) = \sum_{j=0}^{m} a_j(\exp) F_j(A), \quad A \in \mathcal{M}_{n \times n}(C)
\]

(10)

it is well known that

\[
\lim_{m \to \infty} S_m(A)v = \exp(A)v, \quad v \in C^n
\]

(11)

even if \( A \) is not diagonalizable, and this convergence is \textit{asymptotically optimal} with respect to \( K \) (cf. [21] and references therein).

**Remark 2.1.** When \( K \) coincides with the closure of the internal part of an ellipse \( E(d, c, a) \) (symmetric with respect to the real axis, with center \( d \), foci \( d - c \) and \( d + c \), intersections with real axis \( d - a \) and \( d + a \)), it is known that \( \psi(6) \) takes the form

\[
\psi(w) = \gamma w + d + \frac{a - b}{2} \cdot \frac{1}{w} = \gamma \left( w + \frac{d}{\gamma} + \frac{c^2}{4\gamma^2} \frac{1}{w} \right), \quad \gamma = \frac{a + b}{2},
\]

where \( b = \sqrt{a^2 - c^2} \). Hence, the \( m \)-term recurrence of Faber polynomials (7) becomes a three-term recurrence, similar to that of Chebyshev polynomials:

\[
\begin{align*}
F_0(z) &= 1, \\
F_1(z) &= \frac{z}{\gamma} - c_0, \\
F_2(z) &= F_1(z) \cdot F_1(z) - 2c_1, \\
F_j(z) &= F_1(z) \cdot F_{j-1}(z) - c_1F_{j-2}(z), \quad j \geq 3.
\end{align*}
\]

Indeed, in the particular case when \( c \neq 0 \) and \( a \neq b \) (wherefrom \( c^2/(4\gamma^2) \neq 0 \)), it can be proved that the following relation holds

\[
F_j(z) = 2 \left( \frac{c}{\gamma} \right)^j T_j \left( \frac{z - d}{c} \right), \quad j > 0,
\]

where \( T_j \) is the \( j \)-th Chebyshev polynomial (cf. [28]). In the sequel, Faber polynomials and Faber coefficients on an ellipse will be termed (scaled and translated) Chebyshev polynomials and Chebyshev coefficients, respectively.

Notice that we restrict our attention to ellipses symmetric with respect to the real axis since in our application to advection-diffusion equation we deal with real matrices, whose spectrum has the same property.

3. Error bounds for the Faber series approximation to the matrix exponential

In this section we give, on the basis of [19, 21], the convergence estimates for the Faber (Chebyshev) series approximation to the exponential operator. Let \( A \) be diagonalizable with diagonalization matrix \( P \), \( K \) a compact set with capacity \( \gamma \). If \( \sigma(A) \subseteq K \), then

\[
\|\exp(A)v - S_{m-1}(A)v\|_2 \leq k_2(P) \|\exp(\cdot) - S_{m-1}(\cdot)\|_K \|v\|_2,
\]

where \( k_2(P) := \|P\|_2 \cdot \|P^{-1}\|_2 \). If \( \sigma(A) \not\subseteq K \), there exists \( R > \gamma \) so that \( \sigma(A) \subseteq K_R \) (see the previous section for the definition of \( K_R \)) and so

\[
\|\exp(A)v - S_{m-1}(A)v\|_2 \leq k_2(P) \|\exp(\cdot) - S_{m-1}(\cdot)\|_{K_R} \|v\|_2. \tag{14}
\]

When \( K \) is convex, it can be shown (cf. [2, 21]) that, for \( R \geq \gamma \),

\[
\|\exp(\cdot) - S_{m-1}(\cdot)\|_{K_R} \leq \begin{cases} 
8R \left( \frac{4R^2}{\gamma(4R - m)} \right) \left( \frac{m^2}{4R} \right) & \text{if } m \leq 2R \\
4 \left( \frac{\psi_0 \left( \frac{m}{\gamma} \right)}{e \cdot R} \right) & \text{if } m > 2R,
\end{cases} \tag{15}
\]

where

\[
\psi_0(w) := \psi(w) - \gamma w, \quad |w| \geq 1.
\]

Clearly, the estimate is valid for \( R = \gamma \) when \( \sigma(A) \subseteq K = K_\gamma \). Since \( \psi_0 \left( \frac{m}{\gamma} \right) = d + O \left( \frac{1}{m} \right) \) for \( m \geq 2\gamma \), Faber series convergence for the (matrix) exponential is superlinear.
Remark 3.1. It can be proved that the polynomial $S_m$ does not depend on $R$, while it is manifest that the error bound (15) does. This means that the convergence rate is invariant in the family $\{K_R\}_R$. Therefore, a tight bound for this rate is obtained by minimizing the estimate (15) on all $R$ such that $\sigma(A) \subseteq K_R$.

Let us now restrict to the case when the compact set $K$ is $K \equiv \mathcal{E}(d, c, a)$ (cf. Remark 2.1). If $\sigma(A) \subseteq K = \mathcal{E}(d, c, a)$, it can be easily proved that the first of (15) for $R = \gamma$ is dominated by an increasing function with respect to $\gamma$. However, the first of (15) is never of practical interest since, when dealing with advection-diffusion equations with significant advection (with eigenvalues not all real), asymptotic convergence is attained for $m > 2\gamma$. In the case with “small” advection, convergence can be reached when $m < 2\gamma$ but, in this case, the first of (15) is many orders of magnitude larger than the actual error. Bounds given in [29], valid when the spectrum is contained in a real interval, are tighter and can be used instead.

On the other hand, the second of (15), for $R = \gamma$, is itself an increasing function of $\gamma \leq m/2$. Therefore, among ellipses containing $\sigma(A)$ it is convenient to choose that of smallest capacity, which in the sequel will be termed optimal. If $\sigma(A) \not\subseteq K_\gamma$, as just said, convergence still holds, and the tighter error bound correspond to (15) where $R := \inf\{R : \sigma(A) \subseteq K_R\}$ that is to the ellipse with smallest capacity, among the confocal family $\{K_R\}_R$ containing the spectrum.

4. Numerical implementation of the Chebyshev method

The first step consists in approximating the optimal ellipse which contains all the eigenvalues of $A$. To this aim, only the marginal eigenvalues of $A$ are needed. These can be computed by ARPACK (ARnoldi Package †), a collection of FORTRAN77 subroutines based upon an algorithmic variant of the Arnoldi process, called “Implicitly Restarted Arnoldi Method” (cf. [15]). This software is designed to compute a few, say $k$, eigenvalues with user specified features such as those of largest real part or largest magnitude using $n \cdot O(k) + O(k^2)$ storage, where $n$ is the dimension of $A$.

In practice, for advection-diffusion matrices arising from stable centered FD discretizations (grid-Péclet numbers sufficiently smaller than 1, cf. [22]), the convex hull of the spectrum turns out to be very similar to a (possibly degenerating) ellipse; see Figure 1 for an example (here spectra have been computed by the subroutine dgeev of LAPACK).

In view of the shape of the spectrum, we approximate its convex hull by the rectangle constructed using the extremal eigenvalues (i. e. with extremal real and imaginary part) and we select, among the ellipses circumscribing this rectangle, that of smallest capacity (smallest sum of semi-axes). As observed in the previous section, the convergence rate of the corresponding Chebyshev series is driven by the confocal ellipse with smallest capacity containing the spectrum. The latter ellipse is not, in general, the optimal one, but it is expected to have a capacity close to the optimal, as confirmed by our numerical experiments.

To evaluate the extremal eigenvalues (smallest and largest real part, largest imaginary part) we used driver dndrv1 of ARPACK (which does not solve linear systems). Concerning the construction of the ellipse above, denote with $x$, $X$ and $Y$ the smallest real part, the largest

†available from ftp://ftp.caam.rice.edu/pub/software/ARPACK

Figure 1. Eigenvalues distribution of advection-diffusion matrices with $\theta_1 = 60, \theta_2 = 60$ and $n = 1600$ (top) and with $\theta_1 = 40, \theta_2 = 40$ and $n = 1600$ (bottom).
real part and the largest imaginary part of $\sigma(A)$, respectively. A rectangle $R$ of vertices $x + iY$, $X + iY$ and the symmetric to real axis is defined: it is easy to verify that the ellipse $\sigma_R^+$ with center $d = (X + x)/2$ and

$$c^2 = (C^2 + Y^2)(C^2 - Y^2), \quad \gamma = \frac{C^2 \sqrt{C^2 + Y^2} + \sqrt{(CY)^2 + Y^2}}{2}, \quad \left( C := \frac{X - x}{2} \right) \quad (16)$$

is that circumscribed to $R$ with smallest capacity.

Once the ellipse has been determined, Chebyshev coefficients can be evaluated. To this aim, we used a numerical integration scheme based on trapezoidal rule and the fast Fourier transform (FFT) algorithm (cf. [10]). If $K$ is an ellipse, we can choose $R = \gamma$ in (9), whereas if $K$ is a segment of capacity $\gamma'$ we can choose an ellipse of capacity $\gamma'_1 = 1 \times \gamma$ and $R = \gamma'$ (this is possible because an ellipse is a Jordan curve, cf. [7]) and applying the transformation of variables $e^{i2\pi \theta} = w$, we have

$$a_j^{(\text{exp})} = \int_0^1 \exp \left( \left( \gamma + \frac{c^2}{4\gamma} \right) \cos 2\pi \theta + d + i \left( \gamma - \frac{c^2}{4\gamma} \right) \sin 2\pi \theta \right) e^{-ij2\pi \theta} d\theta;$$

using a $(M+1)$-point trapezoidal discretization ($M = 2^N$, $N > 1$) we have, for $0 \leq j \leq M - 1$,

$$a_j^{(M)}(\text{exp}) = \frac{1}{M} \sum_{k=0}^{M-1} \exp \left( \left( \gamma + \frac{c^2}{4\gamma} \right) \cos 2\pi \frac{k}{M} + d + i \left( \gamma - \frac{c^2}{4\gamma} \right) \sin 2\pi \frac{k}{M} \right) e^{-ij2\pi \frac{k}{M}} \quad (17)$$

and so we can use FFT. From error bounds (15), since $2\gamma$ iterations are necessary to reach superlinear convergence, we calculate a priori a safety number of Chebyshev coefficients, say $2N^\ast$, where $N^\ast$ is the smallest integer such that $2N^\ast > \lfloor 4\gamma \rfloor$.

At this point, using (10–12) we could compute in principle the Chebyshev series approximating $\exp(A)v$, stopping the iteration at a suitable index $m$, since theoretical convergence is guaranteed by estimate (15). In practice, however, convergence may not take place when the capacity of the smallest ellipse confocal to that numerically computed, and containing the spectrum, is large. Indeed, as already observed in [24], in this case one is forced to use Chebyshev coefficients affected by serious cancellation errors. There are two main strategies to overcome this problem, both adopted in our implementation: first, computing the Chebyshev coefficients in higher precision and, when this does not suffice, reducing the step-size (fractionizing the capacity), i.e. computing $\exp(A)v$ as

$$\exp(A)v = \exp(\eta_1 A)(\exp(\eta_2 A) \ldots (\exp(\eta_p A) \ldots))v, \quad \sum_{i=1}^{p} \eta_i = 1. \quad (18)$$

We now give and make some comments on our implementation of Chebyshev series approximation of $\exp(A)v$.
Algorithm 4.1

1. INPUT: $A$, $v$, $tol$, $maxitr$, $x$, $X$, $Y$
2. Compute $c^2, \gamma$ by (16)
3. IF $Y/C < 10^{-3}$ OR $C/Y < 10^{-3}$ THEN $\gamma := \gamma \times 1.1$
4. $E := \frac{1}{\gamma}A - \frac{d}{\gamma}I$
5. IF $(4\gamma \geq maxitr)$ THEN
6. $nstep := \left\lfloor \frac{d}{\gamma}maxitr \right\rfloor$
7. $\gamma := \gamma \cdot nstep$, $c := c \cdot nstep$
8. ELSE
9. $nstep := 1$
10. END IF
11. $istep := 0$, $\beta := \|v\|_2$, $conv := $ FALSE
12. DO WHILE (NOT conv)
13. $m_{max} := \lceil 4\gamma \rceil$, $M := \max\{128, 2^{\lceil \log_2 (4\gamma) \rceil}\}$
14. Compute the first $M$ Chebyshev coefficients by (17)
15. $conv :=$ TRUE
16. DO WHILE (($istep < nstep$) AND (conv))
17. $f_0 := v$, $s_0 := a_0 f_0$
18. $f_1 := Ev$, $s_1 := s_0 + a_1 f_1$
19. $r_1 := s_1 - s_0$, $f_0 := 2f_0$, $k := 1$
20. DO WHILE (($\|r_k\|_2 > \|r_{k-1}\|_2 \cdot \|v\|_2$) OR ($k < \gamma$)) AND (conv))
21. $k := k + 1$
22. $f_k := Ef_{k-1} - \frac{c^2}{4\gamma} f_{k-2}$, $s_k := s_{k-1} + a_k f_k$
23. $r_k := s_k - s_{k-1}$
24. IF (($k > 2\gamma$ AND $\|r_k\|_2 > \|r_{k-1}\|_2$) OR ($k = m_{max}$)) conv := FALSE
25. END DO
26. IF (conv) THEN
27. $v := s_k$, $istep := istep + 1$
28. ELSE
29. $nstep := 2 \times (nstep - istep)$, $istep := 0$
30. $\gamma := \frac{\gamma}{2}$, $c := \frac{c}{2}$
31. END IF
32. END DO
33. END DO
34. OUTPUT: $s_k \approx e^A v$

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4.1. Comments on Algorithm 4.1

1. We assume that a suitable approximation of the extremal eigenvalues of $A$, say $x$, $X$, $Y$, is known. This will be useful in the application to advection-diffusion equations where $A = \tau B$ for different $\tau$-values and the spectrum of $A$ can be estimated by that of $B$, which is computed only once. The parameter $\text{maxitr}$ is the maximum degree of the truncated Chebyshev series, and $\text{tol}$ is the exit tolerance relative to $\|v\|_2$.

3. Manages the case in which the ellipse degenerates into a segment.

4. Also this matrix can be computed, as the extremal eigenvalues, only once, being independent of any scaling of $A$, and it is displayed for sake of clarity.

5.-10. Fractionizes the capacity if it is too large, to guarantee convergence using less than $\text{maxitr}$ Chebyshev iterations; $\text{nstep}$ is the number $p$ of steps in (18).

11. $\text{conv} = \text{false}$ when convergence is not reached, and $\gamma$ needs to be furtherly fractionized.

12.-33. Main loop: manages possible fractionizing.

13. $m_{\text{max}} = 4\gamma$ is the maximum number of allowed Chebyshev iterations taking into account that after $2\gamma$ iterations superlinearity should begin (cf. (15)) as confirmed by our numerical results. A minimum value of $M$ must be given (heuristically set to 128) in order to guarantee accuracy of the quadrature rule.

14. Chebyshev coefficients are computed using quadruple precision arithmetics.

16.-32. First level loop: implements the factorization (18).

17.-19. Initialization of the Chebyshev recurrence and series.

20.-25. Second level loop: computes a single exponential operator on the matrix with the (possibly) fractionized capacity. Stopping criteria: relative “residual” (last computed term of the Chebyshev expansion) below the tolerance, provided that a sufficient number of iterations have been done to avoid erroneous convergence.

24. If the residual norm is not decreasing after $2\gamma$ iterations (expected starting of superlinearity, cf. (15) with $R = \gamma$), or the maximum number of iteration $m_{\text{max}}$ has been reached, then convergence has not occurred and we need to fractionize $\gamma$.

26.-31. If $\text{conv} = \text{true}$, proceeds with the stepwise computation of the exponential with unchanged $\gamma$. Otherwise, $\gamma$ is halved and the number of remaining steps is consequently doubled.

5. Numerical results

In this Section we consider the spatial discretization of the linear advection-diffusion equation (1) by five point Finite Differences with constant step-size $h$ on $\Omega = (0,1) \times (0,1)$, which yields the large system of ODEs (2), where

$$B = \frac{1}{h^2} \text{tridiag}(D_+, G, D_-) \in \mathcal{M}_{n \times n}(\mathbb{R}), \quad h = \frac{1}{\nu + 1}, \quad \nu^2 = n$$

(19)

and

$$D_+ = \left(1 + \frac{\theta_2 h}{2}\right) I_\nu, \quad G = \text{tridiag}\left(1 + \frac{\theta_1 h}{2}, -4, 1 - \frac{\theta_1 h}{2}\right), \quad D_- = \left(1 - \frac{\theta_2 h}{2}\right) I_\nu.$$

Below, we compare the Chebyshev method as implemented in the previous Section with the well-known Krylov method, as implemented by Saad in the EXPPRO routine [23, 9], in
the evaluation of $\exp(\Delta t B)\mathbf{v}$ (where $\mathbf{v} := (1,\ldots,1)^T$) for different values of the timestep $\Delta t$, at a given tolerance $tol = 10^{-8}$ (relative to $\|\mathbf{v}\|_2$). We fixed the spatial discretization corresponding to $h = 1/101$ (leading to $n = 10000$) and select three values of the pair $(\theta_1,\theta_2)$, namely $(100,100)$, $(0,50)$, $(50,50)$, which lead to discretization matrices with different spectral features. In all cases the convex hull of the spectrum of $B$ resembles an ellipse (cf. Figure 1) with $C/Y \approx 6.4, 137, 1310$ respectively.

For each test case we report the results of the two methods for five values of $\Delta t$. The largest $\Delta t$ has been selected, for each test case, in order to yield a nearly steady-state solution (the norm of the solution is approximately four orders of magnitude smaller than the norm of $\mathbf{v}$). The performance of both methods is evaluated in terms of ‘equivalent’ scalar products (e.s.p.), that is operations costing $2n$ flops (a mat-vect product counts as e.s.p. with $l$ the average number of nonzeros per row — in our case $l = 5$) and in terms of total CPU seconds on a 600Mhz Alpha station. We observe that the complexity measure in terms of e.s.p. takes into account only the bulk of both methods, in the sense that the Chebyshev “preprocessing” stage (estimate of the spectrum and computation of the series coefficients) and the Krylov “postprocessing” stage (evaluation of the Hessenberg matrix exponential) are not considered.

As for the Krylov implementation, the maximum dimension $m$ of the Krylov subspace is an input parameter. The input timestep $\Delta t$ is automatically subdivided, in view of (4), into a number $p$ of smallest time intervals $\tau_i$, such that $\exp(\tau_i B)$ times a vector can be evaluated within a number of iterations less or equal than $m$. We employed values of $m$ ranging from 5 to 50 since the optimal value of $m$ is not known a priori and it is problem dependent. In the tables we also report the number $p$ of substeps in which each step is subdivided, the number $m \times p$ of mat-vect products (total number of iterations).

Concerning the Chebyshev method, we used Algorithm 4.1 with $A = \Delta t B$ and maxitr = 500. The tables report the results obtained with three different approximations of the spectrum corresponding to

$$ncv := \text{number of ARPACK iterations}$$

equal to 3, 5 or 10 for the approximation of each extremal eigenvalue, and with the ‘true’ values of the extremal eigenvalues (that is, computed by ARPACK with a relative tolerance of $10^{-6}$), in order to estimate the impact of possible inaccuracy at this stage. Note that, once the extremal eigenvalues of $B$ (say $x_B, X_B, Y_B$) have been approximated, we can use Algorithm 4.1 for every choice of $\Delta t$ after setting: $x = \Delta t x_B, X = \Delta t X_B$ and $Y = \Delta t Y_B$. In the tables we give the number $p$ of substeps, and the number of total mat-vect products (number of iterations) as a sum of those necessary to assess the value of $\gamma$ which guarantees convergence, plus those needed for the actual computation of the right hand side of (4). The total CPU time in seconds and the partial times for eigenvalue approximation, Faber coefficients evaluation, series evaluation, respectively, are also reported.

5.1. Comments on the numerical results

The experimental results, collected in the Tables I–VIII, show that the Chebyshev method with a rough approximation of extremal eigenvalues ($ncv = 10$) performs better than the Krylov one with the optimal choice $m$ of the dimension of the subspace. In Table I we show the evolution of the solution norm for different timesteps $\Delta t$ up to a nearly steady-state for the three test problems.

Tables II, III and IV refer to the test case with $\theta_1 = \theta_2 = 100$. From Table II we note
that the Krylov method is very sensitive to the choice of \( m \), which governs the fractionizing of the timestep and hence the number of substeps \( p \). In particular, when \( m \) is too small, the complexity is affected by the extremely large number \( p \) of substeps, while for large \( m \) the algorithm is penalized by the quadratic complexity (in \( m \)) of the single steps. In practice the ‘optimal’ value of \( m \) ranges between 10 and 20 for this problem. We notice that large values of \( m \) are in general not recommended due to memory occupancy.

In Table III we present the results of the Chebyshev-Algorithm 4.1 for the various timesteps. The timestep is subdivided into \( p \) substeps if:

- the approximation of the eigenvalues is good but the capacity of the ellipse is too large to
guarantee convergence within the allowed maximum number of iterations (the capacity and hence the timestep is fractionized a priori, cf. steps 5–10 of Algorithm 4.1). This occurs, for example, with $\Delta t = 0.01, 0.012$ and $\text{ncv} = 10$ where $p = 2$ and $p = 3$, respectively.

- The extremely rough approximation of the spectrum leads to non-decreasing error estimate after $2\gamma$ iterations, so that the capacity is halved (step 24 of the Algorithm 4.1). This occurs, for instance, with $\Delta t = 0.01$ and $\text{ncv} = 3$, where the timestep is halved 5 times leading to $p = 32$.

The cost of the extremal eigenvalue computation is displayed as the first of the partial CPU seconds. For the two smallest timesteps a very rough approximation of the spectrum (with 3

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<th>e.s.p.</th>
<th>partial CPU s.</th>
<th>CPU s.</th>
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Table III. Results of the Chebyshev method for the test case with $\theta_1 = \theta_2 = 100$.

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<th>partial CPU s.</th>
<th>CPU s.</th>
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Table IV. As in the previous Table with $\text{maxitr} = 1000$. 
Krylov-EXPRO ($\theta_1 = 0$, $\theta_2 = 50$), $tol = 10^{-8}$

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Table V. Results of the Krylov method for the test case with $\theta_1 = 0$, $\theta_2 = 50$

or 5 ARPACK iterations, respectively) gives convergence comparable to that obtained with the ‘true’ eigenvalues. For larger timesteps, 10 ARPACK iterations produce a satisfactory approximation of the spectrum together with a computational cost which is at most the 30% of the whole algorithm. In any case, Chebyshev with $n_{cv} = 10$ for the extremal eigenvalues is more efficient than the Krylov method for all the timesteps. Comparison of Chebyshev with ‘optimal’ Krylov shows the following speedups (ratios between total CPU times): 1.0, 1.2, 2.7, 3.2 and 3.2, from the smallest to the largest $\Delta t$. It is worth noticing the apparently strange feature that the number of mat-vect products using the ‘true’ extremal eigenvalues is larger than using an approximation of the spectrum with $n_{cv} = 10$. There are two reasons for this behavior: one is that the ellipse $E_1$, constructed from the ‘true’ extremal eigenvalues (solid line in Figure 2), has a much larger capacity than the one ($E_2$) constructed using the approximated spectrum (dashed line), thus producing a priori a larger number of substeps for $\Delta t \geq 0.005$ (cf. Table III). The second reason is that the optimal ellipse (dot-dashed in Figure 2) among those confocal to $E_1$, has a larger capacity than the optimal ellipse (dotted line) among those confocal to $E_2$, and then convergence is slower in the first case. This behavior is manifest in Table IV where $maxitr$ has been increased (= 1000) in order to force the method using the exact eigenvalues and that with $n_{cv}=10$ to fractionize a priori (cf. step 5 of Algorithm 4.1) with the same number of substeps $p = 2$: indeed, the number of iterations per step is smaller using the approximated eigenvalues.

The computational results for the second test case are summarized in Tables V and VI.
EFFICIENT APPROXIMATION OF THE EXPONENTIAL OPERATOR

Table VI. Results of the Chebyshev method for the test case with $\theta_1 = 0$, $\theta_2 = 50$

<table>
<thead>
<tr>
<th>$\Delta t$</th>
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<th>mat-vect</th>
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From Table V we see that, as in the previous problem, the optimal $m$-value for the Krylov method lies in the interval $[10, 20]$. Comparing Tables V and VI, we see again that Chebyshev with $\text{ncv} = 10$ is better than 'optimal' Krylov, with the exception of the smallest $\Delta t$. The speedups are: 0.7, 1.3, 2.1, 2.2 and 1.9, respectively.

The results for the third test case are reported in Tables VII and VIII. Also here Chebyshev with $\text{ncv} = 10$ is faster than optimal Krylov, except for $\Delta t = 5 \cdot 10^{-4}$. The speedups are 0.7, 1.2, 2.3, 2.9 and 2.9, respectively.

As observed before, the cost of extremal eigenvalue approximation for the Chebyshev method can be a large percentage of the computation for small timesteps. However, if we have to compute $\exp(\Delta t B)v$ several times, for different $v$, and this is the case when an accurate timestepping is performed, the eigenvalue computation can be carried on only once and for all and becomes completely negligible.

REFERENCES

L. BERGAMASCHI, M. CALIARI AND M. VIANELLO

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Table VII. Results of the Krylov method for the test case with $\theta_1 = \theta_2 = 50$

Table VIII. Results of the Chebyshev method for the test case with \( \theta_1 = \theta_2 = 50 \)

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28. G. Starke and R. S. Varga, *A hybrid Arnoldi-Faber iterative method for nonsymmetric systems of
Figure 2. ‘True’ and approximated extremal eigenvalues and optimal ellipses for the test case with $\theta_1 = \theta_2 = 100$ and $n = 10000$