

Vortex dynamics by geometric and structural complexity analysis

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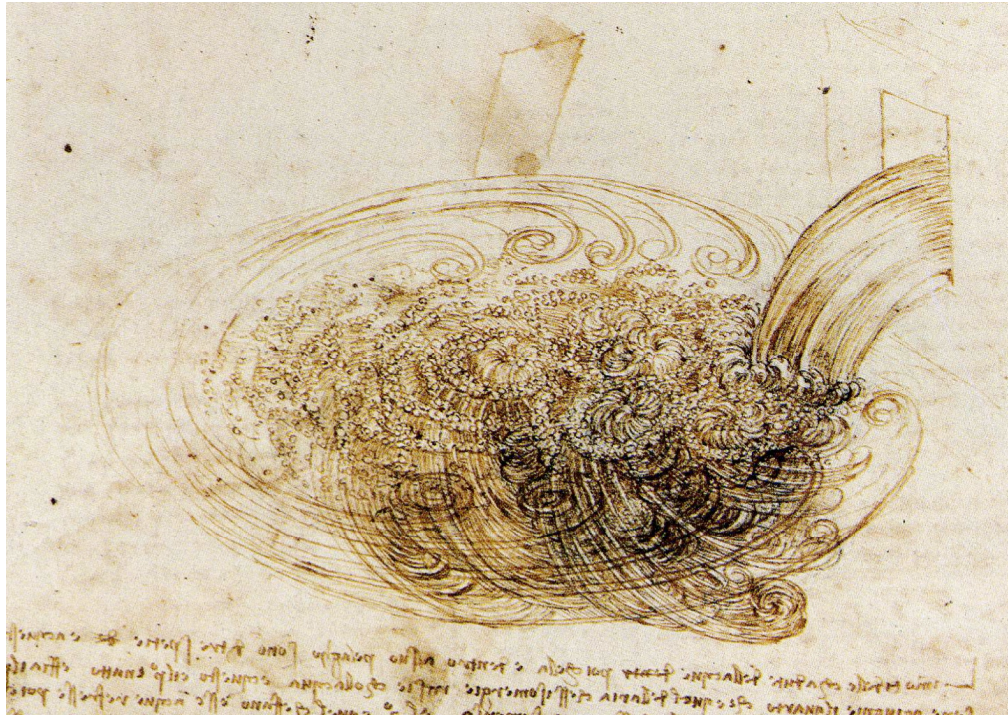
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Aims

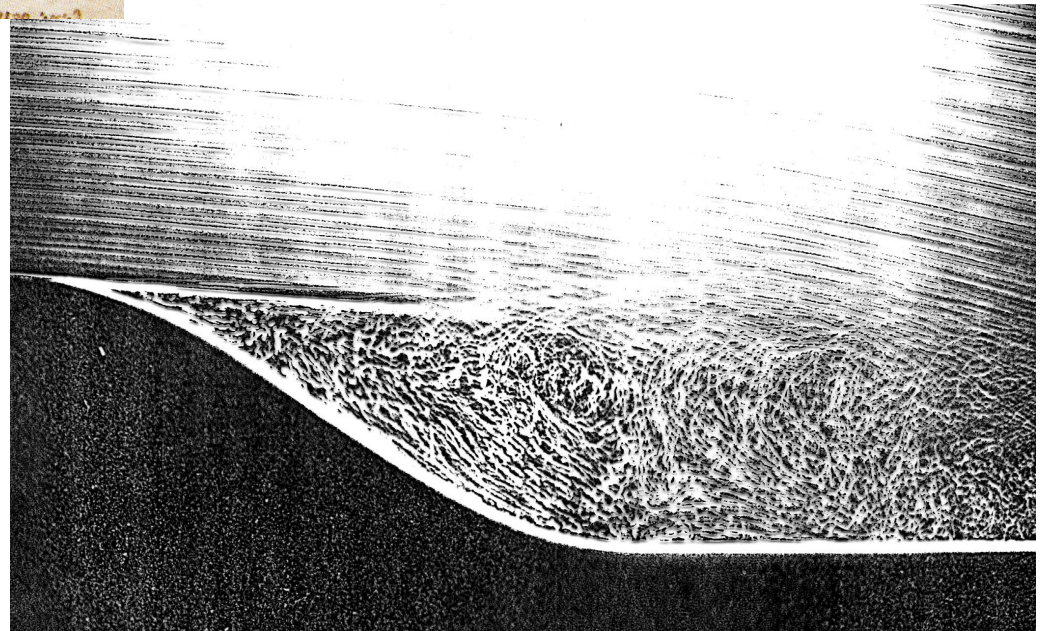
- **Theoretical goals:**
 - describe and classify complex morphologies;
 - study possible relationships between energy and complexity;
 - understand and predict energy localization and transfer;
- **Applications:**
 - implement new visiometric tools and diagnostics;
 - develop real-time energy analysis of dynamical processes.

Coherent structures

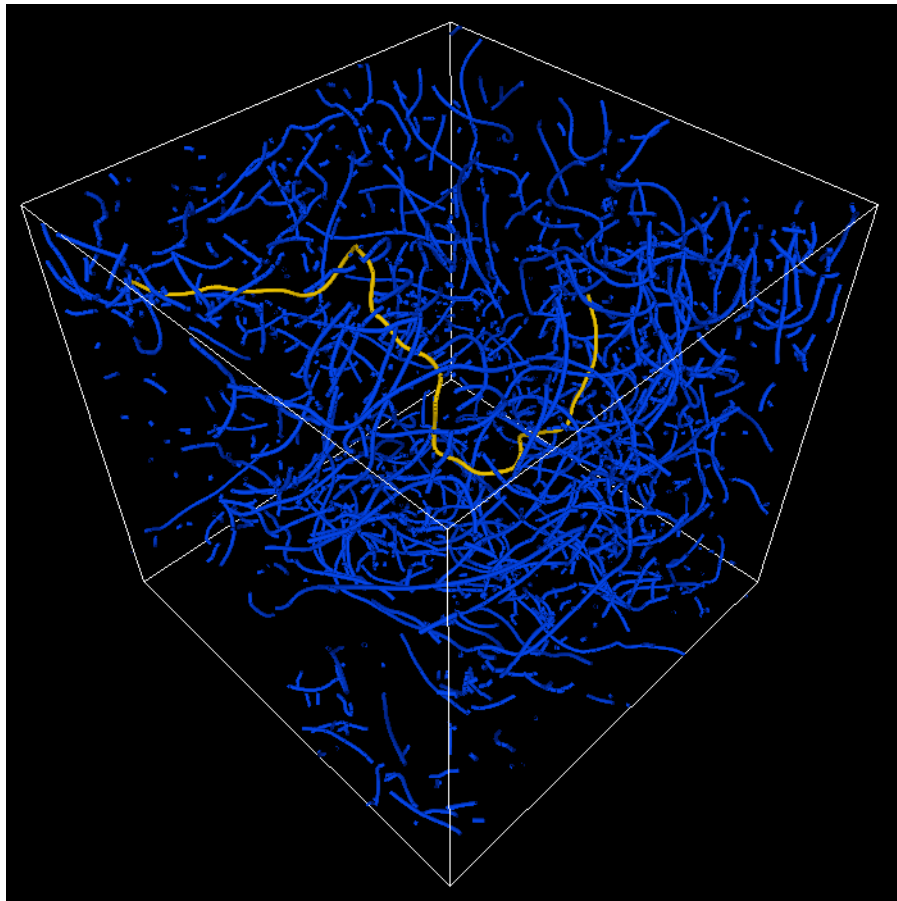


Leonardo da Vinci
(Water studies, 1506)

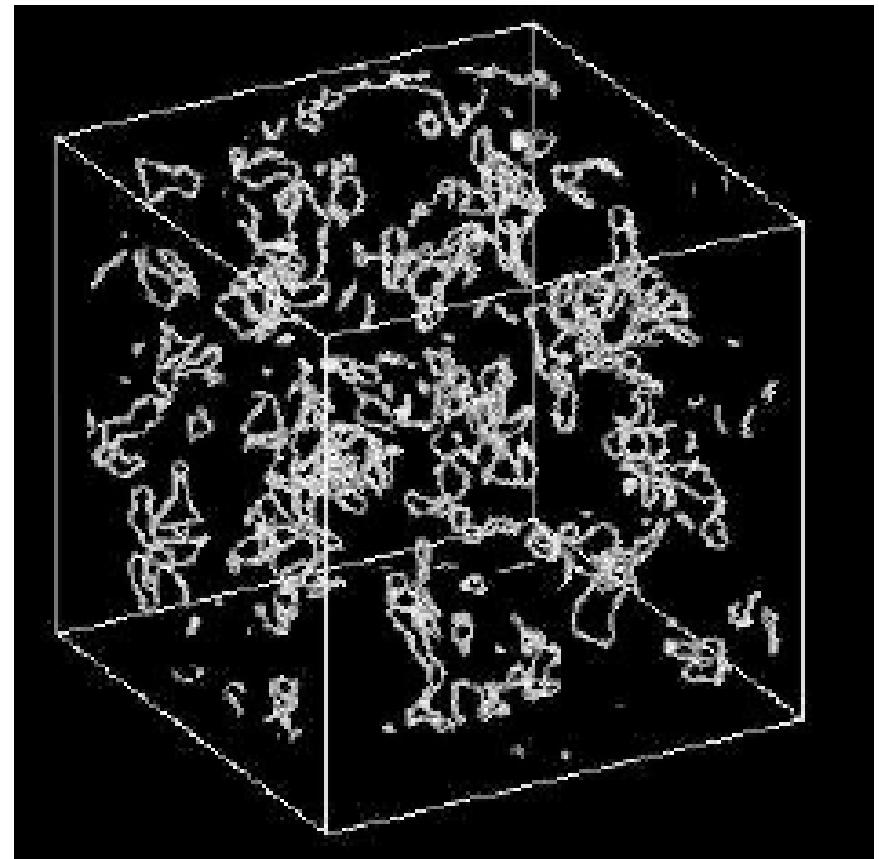
Werlé, ONERA, 1974
(Van Dyke, 1982)



Vorticity localization in classical and quantum fluids



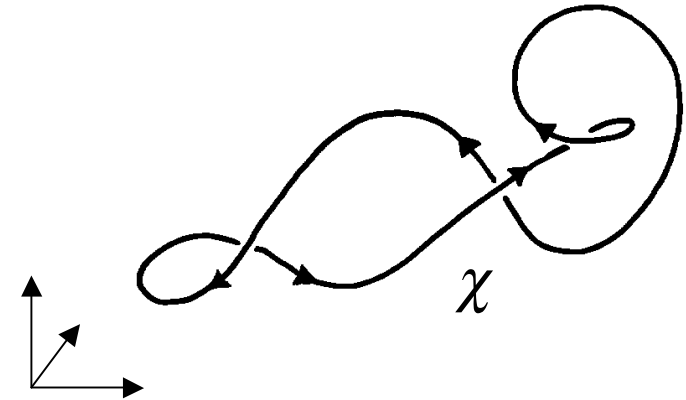
Kida et al.
(Toki-Kyoto, 2002)



Miyazaki et al.
(Physica D, 2009)

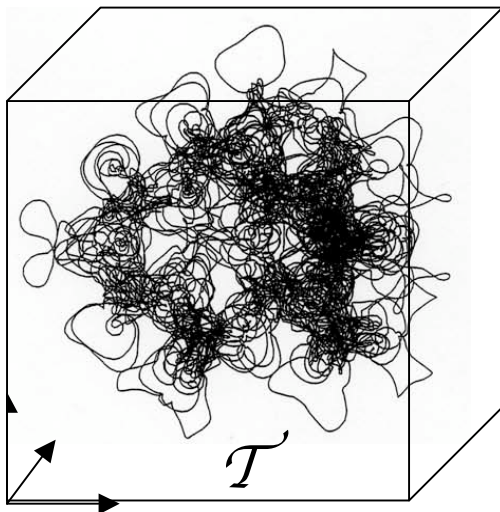
Geometric approach to vortex filament motion

- **Localized induction approximation (LIA) and intrinsic kinematics**
- **Solitons and integrable vortex dynamics**
- **Stationary solutions**
- **Topological properties and fluid invariants**



Structural complexity analysis of vortex dynamics

- **Geometry and topology of fluid flows**
- **Measures of morphological complexity**
- **Energy/complexity relations**
- **Shapefinders and eigenvalue analysis**
- **Dynamical properties in terms of graph analysis**



Localized induction approximation (LIA)

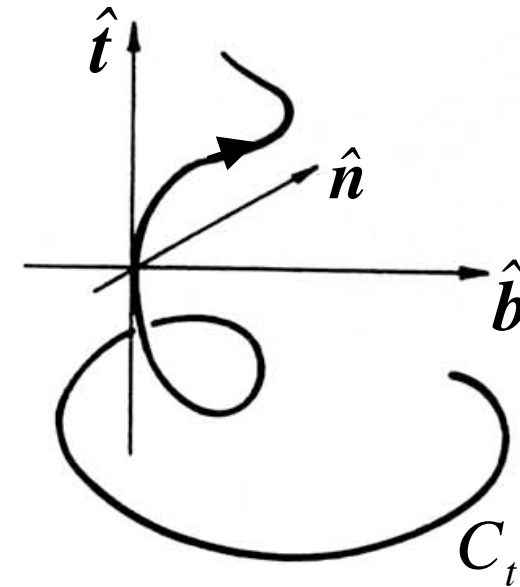
$$\left. \begin{array}{l} \text{homogeneous} \\ \text{incompressible} \\ \text{inviscid} \end{array} \right\} \text{fluid in } \mathbb{R}^3 : \quad u = u(X, t) \quad \left\{ \begin{array}{l} \omega = \nabla \times u \\ \nabla \cdot u = 0 \quad \text{in } \mathbb{R}^3 \\ u = 0 \quad \text{as } X \rightarrow \infty \end{array} \right.$$

- Space curve C_t , given by:

$$C_t : \left\{ \begin{array}{l} X(s, t) := X_t(s) \in C^v \\ s \in [0, L] \rightarrow \mathbb{R}^3 \end{array} \right.$$

- Intrinsic reference on C_t , given by:
(Frenet frame)

$$\hat{t} := X'(s, t) = \frac{\partial X}{\partial s}, \quad \{\hat{t}, \hat{n}, \hat{b}\}$$



- Vortex line on C_t : $\omega = \omega_0 \hat{t}$, $\omega_0 = \text{constant}$

$$\left. \begin{array}{l} \text{asymptotic theory} \\ (R/a \gg 1) \\ \text{no self-intersections} \end{array} \right\} \text{LIA:} \quad \dot{X}(s, t) \equiv \frac{\partial X}{\partial t} \approx \underline{X' \times X''} = c \hat{b} \equiv u_{\text{LIA}}$$

(Da Rios, 1906; Hama & Arms, 1961)

Intrinsic equations under LIA and NLSE

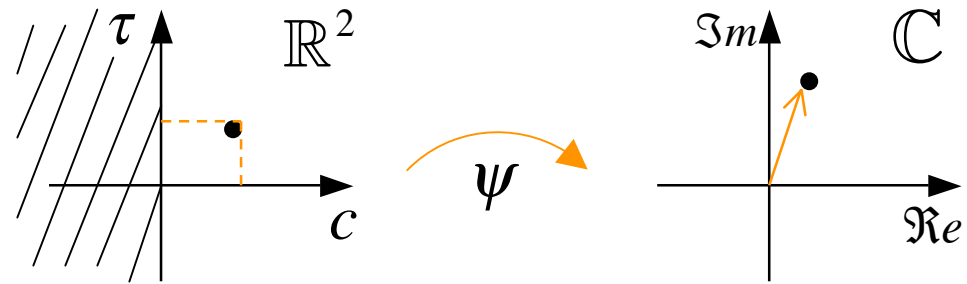
- Intrinsic description:** $\mathbf{u} = (u_t, u_n, u_b)$, $c(s, t)$ **curvature**, $\tau(s, t)$ **torsion**;

under LIA: $\mathbf{u}_{\text{LIA}} = c\hat{\mathbf{b}}$, $u_t = u_n = 0$, $u_b = c$.

$$C_t := \varphi_t(C) \quad \begin{cases} \dot{c} = -(c\tau)' - c'\tau \\ \dot{\tau} = \left(\frac{c'' - c\tau^2}{c} \right)' + c'c \end{cases} \quad \begin{matrix} \text{[Da Rios, 1906]} \\ \text{[Betchov, 1965]} \end{matrix}$$

- NLSE via Madelung transform:**

$$\psi(s, t) = c(s, t) e^{i \int_0^s \tau(\xi, t) d\xi}$$

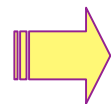


NLSE: $i\psi + \psi'' + \frac{1}{2}|\psi|^2\psi = 0$

(Hasimoto, 1972)

- Da Rios-Betchov eqs. from NLSE:**

by taking log-derivative of $\psi(s, t)$



$$C_t := \varphi_t(C) \quad \begin{cases} \dot{c} = f_c(s, t) \\ \dot{\tau} = f_\tau(s, t) \end{cases}$$

Higher-order LIAs and integrable geometric dynamics

- Higher-order effects:

$$\mathbf{u} = (\alpha + \beta s)c\hat{\mathbf{b}} + (\gamma + \delta s)\hat{\mathbf{t}} + \mu\left(\frac{1}{2}c^2\hat{\mathbf{t}} + c'\hat{\mathbf{n}} + c\tau\hat{\mathbf{b}}\right)$$

β, δ : inhomogeneities
(Lakshmanan & Ganesan, 1985)

γ : non-linear stretching
(Onuki, 1985)

μ : axial flow and vorticity
(Fukumoto & Miyazaki, 1991)

- Higher-order LIAs:

$$\mathbf{u}^{(0)} = \mathbf{u}_{\text{LIA}}$$

$$\mathbf{u}^{(j+1)} = \sum_{v=1}^{j+1} \mathcal{F}\left(\frac{\partial^v \mathbf{u}^{(0)}}{\partial s^v}\right) = \mathcal{R}(\mathbf{u}^{(0)})$$

(Langer & Perline, 1991)

$$\dot{\psi} = \mathcal{P}_\psi(u_n e^{i\Theta}) + \mathcal{Q}_\psi(u_b e^{i\Theta})$$

(Nakayama et al., 1992)

- Integrable geometric dynamics:

$$\mathbf{u} = (u_t, u_n, u_b) \quad \Rightarrow \quad C_t := \varphi_t(C) \quad \begin{cases} \dot{s} = f_s(t) \\ \dot{c} = f_c(s, t) \\ \dot{\tau} = f_\tau(s, t) \end{cases} \quad \left[\begin{array}{l} \text{Germano, 1983} \\ \text{Ricca, 1991} \end{array} \right]$$

Conserved quantities from integrable dynamics

$$A = \oint X \times X' ds = \text{const.} \quad \left[\text{Arms \& Hama, 1965} \right]; \quad L = \oint ds = \text{const.} \quad \left[\text{Marsden \& Weinstein, 1983} \right]$$

- Geometric invariants:

$$\mathcal{P}(L) = \sum_{j=0}^{\infty} \mathcal{I}_j L^j$$

(Langer & Perline, 1991; Ricca, 1991)

$$\mathcal{I}_0 = \oint \tau ds$$

$$\mathcal{I}_1 = \oint c^2 ds$$

$$\mathcal{I}_2 = \oint c^2 \tau ds$$

$$\mathcal{I}_3 = \oint \left(\frac{c^4}{4} - c'^2 - c^2 \tau^2 \right) ds$$

... ..

- Physical conserved quantities:

$\left[\text{Fukumoto, 1987} \right]$
 $\left[\text{Ricca, 1992} \right]$

kinetic energy:

$$\mathcal{K} = \frac{1}{2} \int \|u\|^2 d^3 X$$

pseudo-helicity:

$$\mathcal{H} = \int u \cdot (\nabla \times u) d^3 X$$

enstrophy:

$$\mathcal{E} = \int \omega d^3 X$$

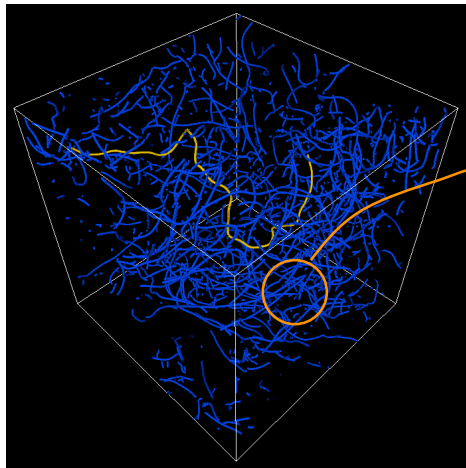
linear momentum:

$$P = \frac{1}{2} \int (X \times \omega) d^3 X$$

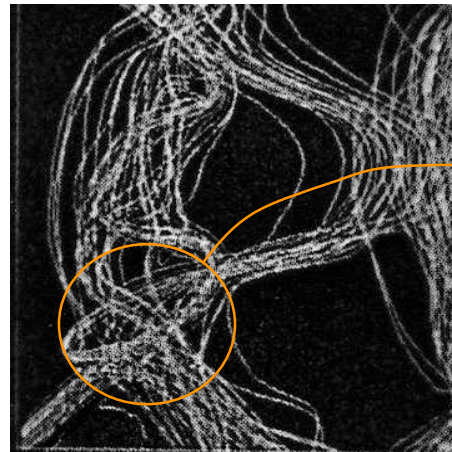
angular momentum:

$$M = \frac{1}{3} \int X \times (X \times \omega) d^3 X$$

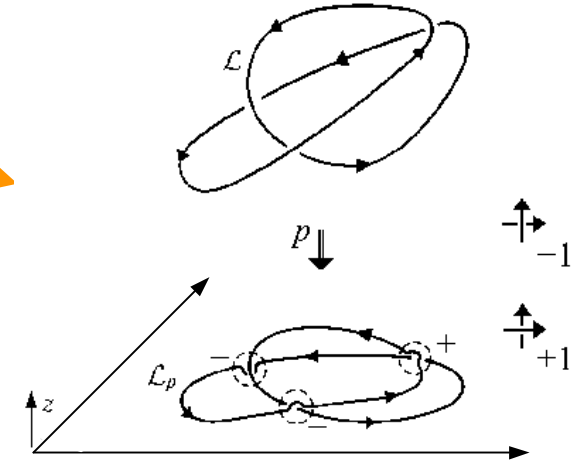
Structural complexity analysis of 3D vortex tangles



**experiment
or simulation**



domain extraction



**projected diagram
analysis**

● Measures of structural complexity:

Algebraic information:

- average crossing number

Topological information:

- minimal crossing number
- linking numbers
- topological changes

Geometric information:

- tropicity directions
- coiling, writhing
- alignment
- signed area

Dynamical systems analysis:

- topological entropy
- eigenvalue analysis

Energy and helicity of a vortex tangle

- **Kinetic energy:** $E(\mathcal{T}') \equiv \frac{1}{2} \int_{\mathcal{T}'} \|\mathbf{u}\|^2 d^3X = \int_{\mathcal{T}'} \mathbf{u} \cdot (\mathbf{X} \times \boldsymbol{\omega}) d^3X$;

from **Lamb (1932)**, we have:
$$E(\mathcal{T}') \approx \frac{\Gamma^2}{8\pi} \sum_{ij} \iint_{\chi_i \chi_j} \frac{\hat{\mathbf{t}}_i \cdot \hat{\mathbf{t}}_j}{|\mathbf{X}_i - \mathbf{X}_j|} ds_i ds_j ,$$

with total length of vortex filaments given by
$$L(\mathcal{T}') = \sum_i \int_{\chi_i} \hat{\mathbf{t}}_i ds_i .$$

- If $\boldsymbol{\omega} \equiv \nabla \times \mathbf{u} = \boldsymbol{\omega}_0 \hat{\mathbf{t}}$, then the **kinetic helicity** is given by
(**Moffatt, 1969; Moffatt & Ricca, 1992**):

$$\mathcal{H}(\mathcal{T}') \equiv \int_{\mathcal{T}'} \mathbf{u} \cdot \boldsymbol{\omega} d^3X = \sum_i Lk_i \Gamma_i^2 + 2 \sum_{i \neq j} Lk_{ij} \Gamma_i \Gamma_j$$

- $Lk_i = Lk(\chi_i, \chi_i) = Wr_i + Tw_i$ **Calugareanu-White invariant**

where

- $Lk_{ij} = Lk(\chi_i, \chi_j)$ **Gauss linking number**

- Γ_i **i-th vortex circulation**

Tangle analysis by indented projections

Let $\Pi_i = \Pi(\hat{T}_i)$ be the “indented” $\hat{T}_i$ -projection of the oriented tangle component χ_i ; assign the value $\varepsilon_r = \pm 1$ to each apparent crossing in Π_i .

- writhing:**

$$Wr_i = Wr(\chi_i) = \left\langle \sum_{r \in \chi_i} \varepsilon_r \right\rangle, \quad Wr = Wr(\mathcal{T}) = \left\langle \sum_{r \in \mathcal{T}} \varepsilon_r \right\rangle ;$$

- linking:**

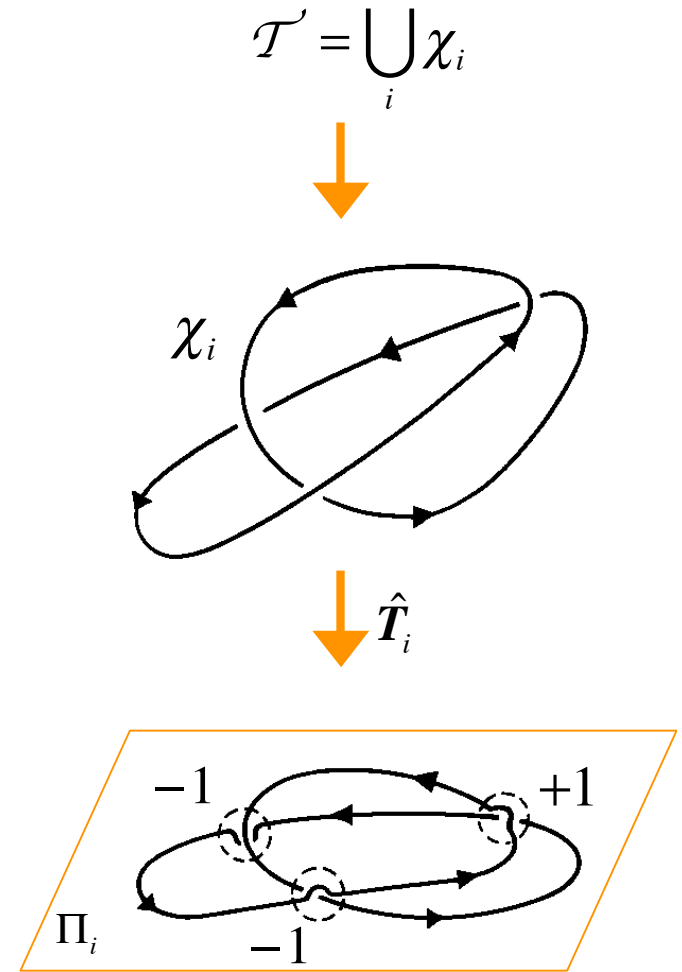
$$Lk_{ij} = Lk(\chi_i, \chi_j) = \frac{1}{2} \sum_{\substack{r \in \chi_i \cap \chi_j \\ i \neq j}} \varepsilon_r, \quad Lk_{tot} = \sum_{\substack{r \in \mathcal{T} \\ i \neq j}} |Lk_{ij}| ;$$

- average crossing number:**

$$\bar{C}_{ij} = \bar{C}(\chi_i, \chi_j) = \left\langle \sum_{r \in \chi_i \cap \chi_j} |\varepsilon_r| \right\rangle, \quad \bar{C} = \sum_{r \in \mathcal{T}} \bar{C}_{ij} ;$$

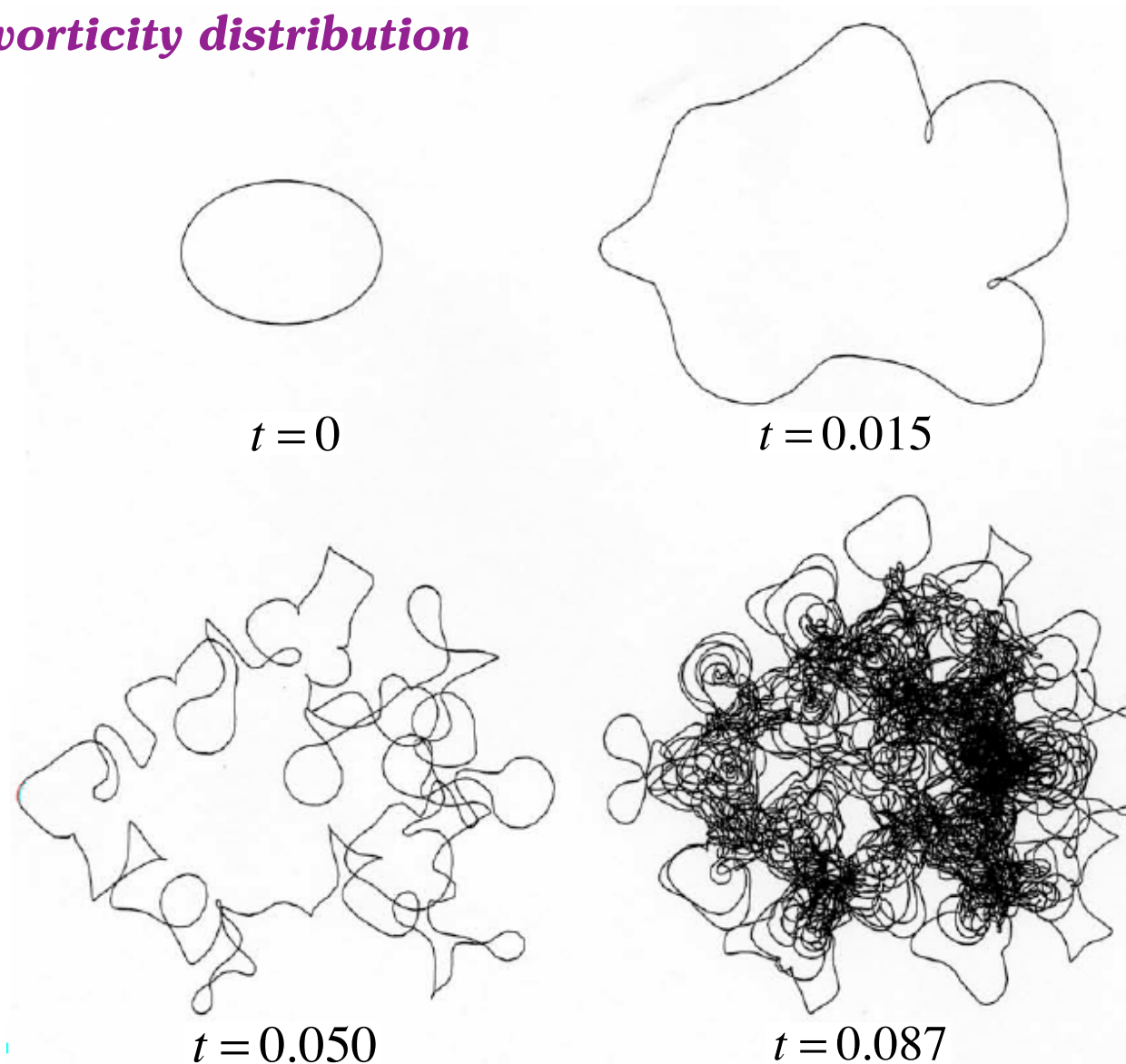
- estimated values:**

$$\bar{Wr}_{\perp} = \left(\sum_{r \in \mathcal{T}} \varepsilon_r \right)_{\perp}, \quad \bar{C}_{\perp} = \left(\sum_{r \in \mathcal{T}} |\varepsilon_r| \right)_{\perp}.$$

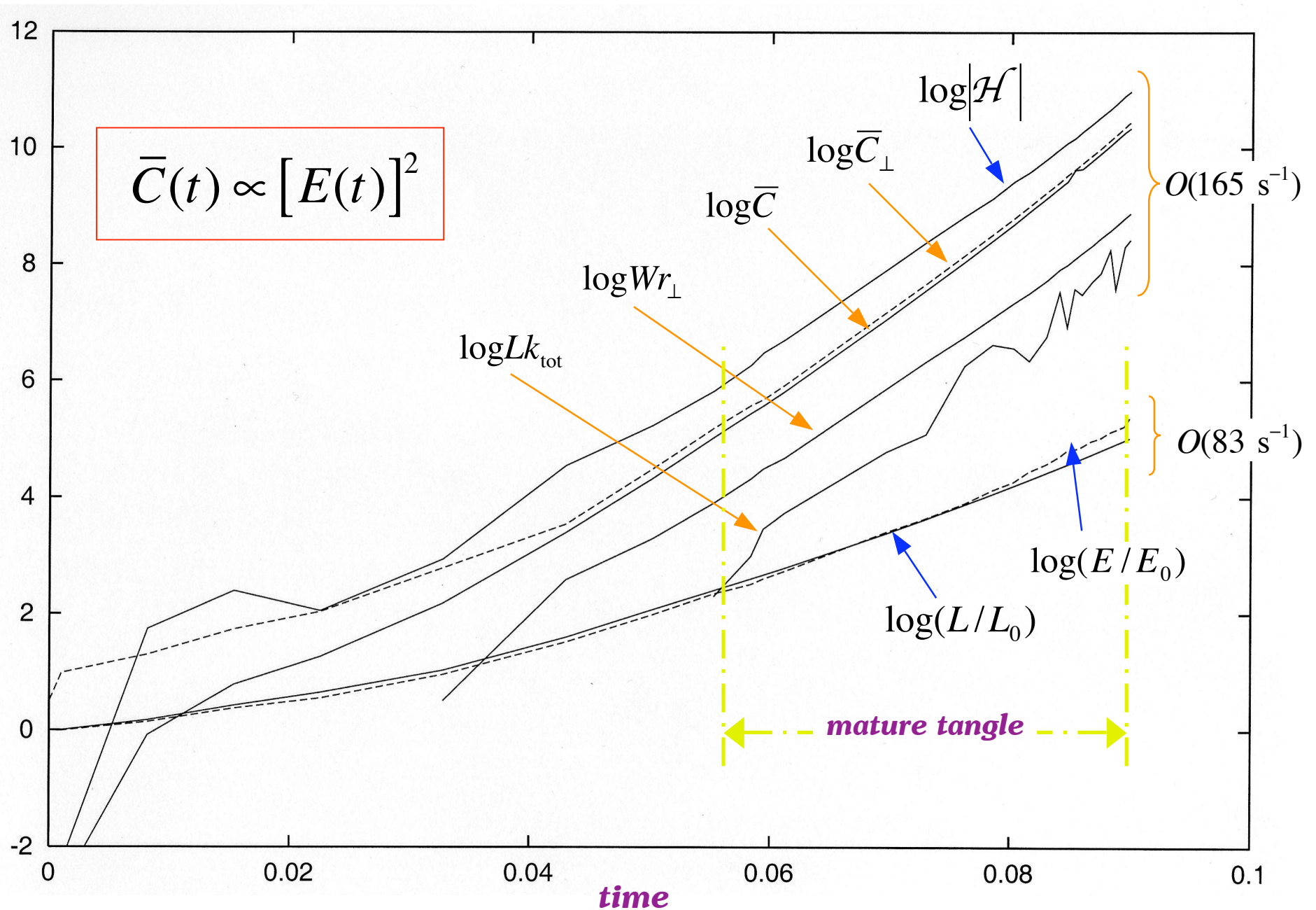


A vortex tangle test case (Barenghi, Ricca & Samuels, 2001)

- ABC-type flow field super-imposed on initial vorticity distribution



Vortex tangle test case: energy-complexity relation



Energy-complexity relation

- $\bar{C}(t) \propto [E(t)]^2$: **since**

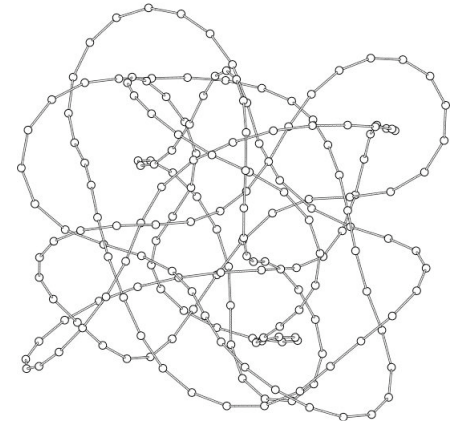
$$\bar{C} = \frac{1}{4\pi} \sum_{ij} \iint_{\chi_i \chi_j} \frac{|(X_i - X_j) \cdot \hat{t}_i \times \hat{t}_j|}{|X_i - X_j|^3} ds_i ds_j ,$$

we have

$$\delta \bar{C} \propto \left(\frac{\delta s}{|X_i - X_j|} \right)^2 , \quad \text{and} \quad \delta(E / E_0) \propto \frac{\delta s}{|X_i - X_j|} ;$$

hence

$$\bar{C} \propto (E / E_0)^2 .$$



Helicity-complexity bounds

If $Lk_{ii} = 0$, $\Gamma_i = \Gamma \quad \forall \chi_i \in \mathcal{T}$, **then we have (Ricca, 2008):**

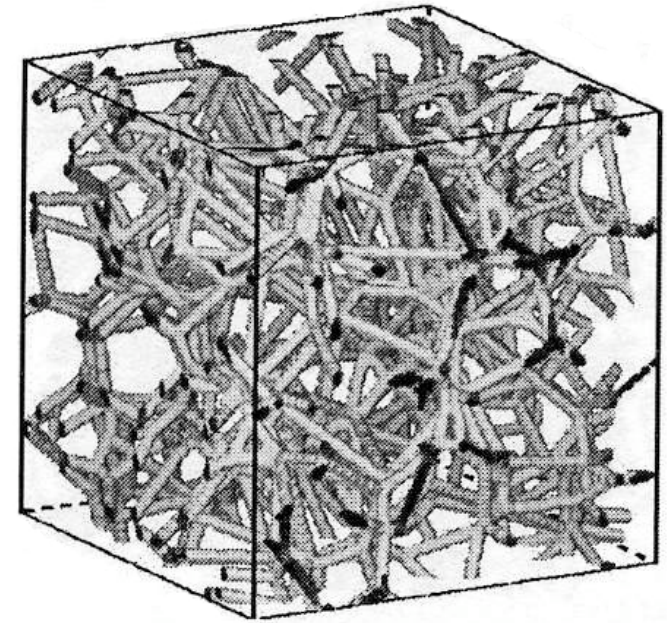
- $|H(t)| \leq 2\Gamma^2 \bar{C}(t) \quad \text{and} \quad \left| \frac{d|H(t)|}{dt} \right| \leq 2\Gamma^2 \frac{d\bar{C}(t)}{dt} .$

Tropicity interpretation of eigenvalues

- **Eigenvalue analysis:** let $\lambda_1 \geq \lambda_2 \geq \lambda_3$.

From (Vilanova et al., 2006), we have:

$$C_f = \frac{\lambda_1 - \lambda_2}{\sum_i \lambda_i} , \quad C_p = \frac{2(\lambda_2 - \lambda_3)}{\sum_i \lambda_i} , \quad C_s = \frac{3\lambda_3}{\sum_i \lambda_i} .$$



- **Shapefinders:** let

$$V = \int_{\mathcal{D}} d^3 X , \quad A = \int_{\partial \mathcal{D}} d^2 X , \quad H_G = \frac{1}{2} \int_{\partial \mathcal{D}} \left(\frac{1}{R_1} + \frac{1}{R_2} \right) d^2 X , \quad \chi_E = \frac{1}{2} \int_{\partial \mathcal{D}} \frac{1}{R_1 R_2} d^2 X .$$

Then

$$L = \frac{H_G}{4\pi\chi_E} , \quad W = \frac{A}{H_G} , \quad T = \frac{3V}{A} ; \quad (\text{Sahni et al., 1998})$$

If $\mathcal{D}$ is convex and $L \geq W \geq T > 0$, we can define

$$C_F = \frac{L - W}{L + W + T} , \quad C_P = \frac{2(W - T)}{L + W + T} , \quad C_S = \frac{3T}{L + W + T} ,$$

so that

$$\lambda_1 \propto L , \quad \lambda_2 \propto W , \quad \lambda_3 \propto T .$$

Interpretation of momenta in terms of projected areas

- **Linear momentum:**

$$\mathbf{P} \equiv \frac{1}{2} \int_{\mathcal{T}} \mathbf{X} \times \boldsymbol{\omega} \, d^3\mathbf{X} = \frac{1}{2} \sum_i \Gamma_i \oint_{\chi_i} \mathbf{X} \times \mathbf{X}' \, ds ,$$

- **Angular momentum:**

$$\mathbf{M} \equiv \frac{1}{3} \int_{\mathcal{T}} \mathbf{X} \times (\mathbf{X} \times \boldsymbol{\omega}) \, d^3\mathbf{X} = \frac{1}{3} \sum_i \Gamma_i \oint_{\chi_i} \mathbf{X} \times (\mathbf{X} \times \mathbf{X}') \, ds .$$

- **Under Euler (and LIA) equations:** $\frac{d\mathbf{P}}{dt} = 0 , \quad \frac{d\mathbf{M}}{dt} = 0 .$

Plane graph by standard projection

$$\Lambda \equiv p(\mathcal{T}) : \mathbb{R}^3 \rightarrow \mathbb{R}^2 \quad \left\{ \begin{array}{l} \Lambda_{xy} \equiv p_z(\mathcal{T}) \\ \mathcal{A}_{xy} \equiv \mathcal{A}(\Lambda_{xy}) \end{array} \right\}_{xy} , \quad \left\{ \dots \right\}_{yz} , \quad \left\{ \dots \right\}_{zx} .$$

Then (Arms & Hama, 1965; Ricca, 1992):

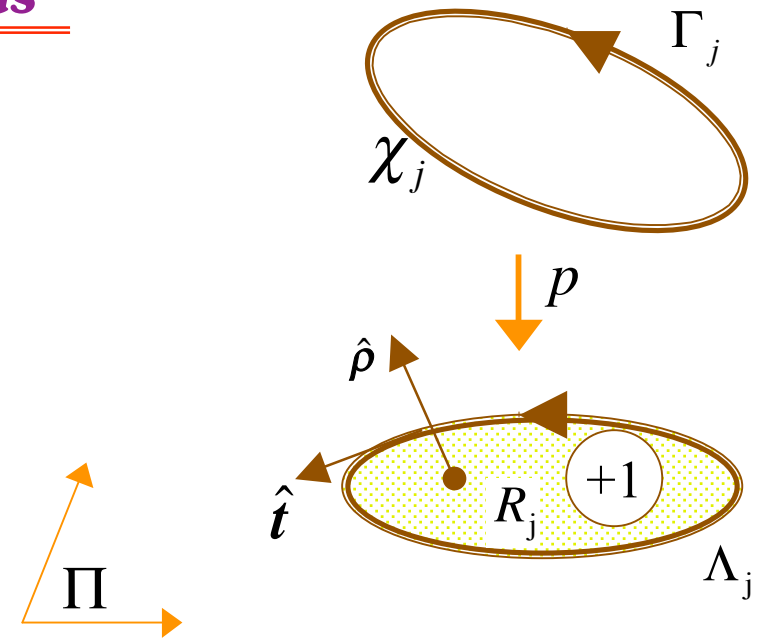
- $P_{xy} = \bar{\Gamma} \mathcal{A}_{xy} , P_{yz} = \bar{\Gamma} \mathcal{A}_{yz} , P_{zx} = \bar{\Gamma} \mathcal{A}_{zx} ;$
- $M_{xy} = \bar{\Gamma} d_z \mathcal{A}_{xy} , M_{yz} = \bar{\Gamma} d_x \mathcal{A}_{yz} , M_{zx} = \bar{\Gamma} d_y \mathcal{A}_{zx} .$

Signed area and weight of projected graphs

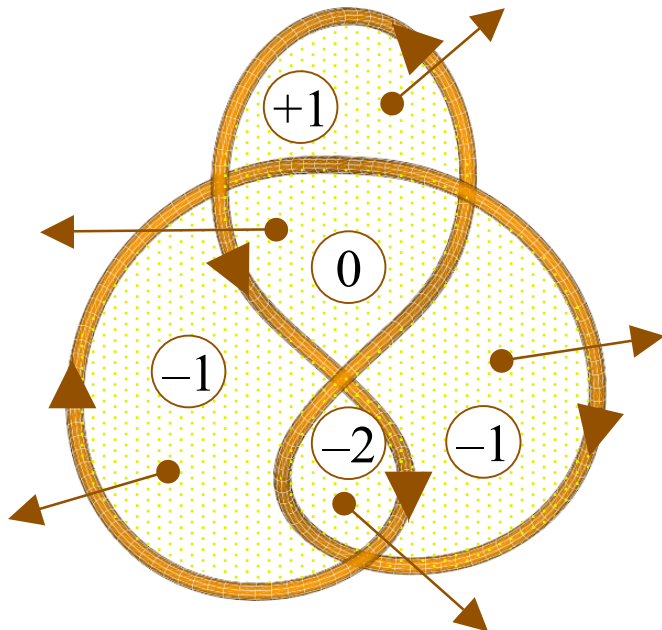
Let $\Lambda_j = p(\chi_j)$ and $R_j \equiv \text{int}(\Lambda_j \cap \Pi)$:

- **Cauchy index** $\mathcal{I}_j = \mathcal{I}(\Lambda_j)$: at each intersection $\hat{\rho} \cap \bar{\Lambda}_j$ assign $\varepsilon_i = \pm 1$ and take

$$\mathcal{I}_j = \mathcal{I}(\Lambda_j) = \sum_{\hat{\rho} \cap \bar{\Lambda}_j} \varepsilon_i .$$



Example



- **Signed area:** $\mathcal{A}_j(\Lambda) = \sum_j \mathcal{I}_j \mathcal{A}_j(R_j)$

- **Weighted circulation:** $\gamma_j = \frac{\sum_k \Gamma_k L_{k,j}}{L_j}$

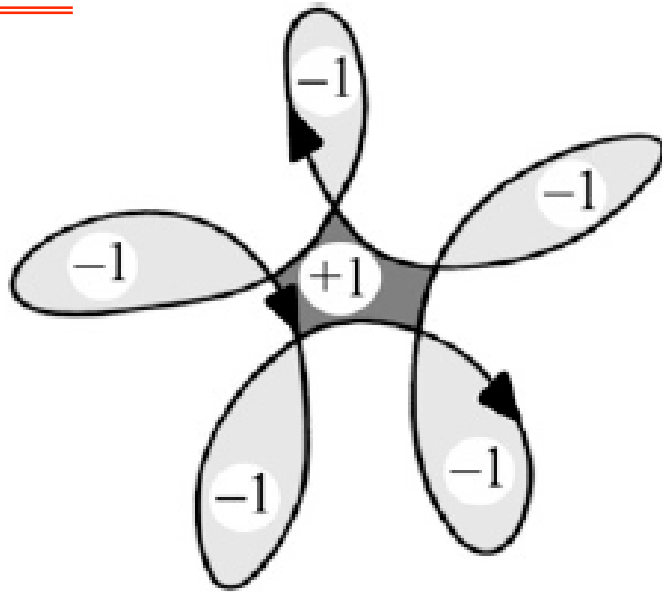
Momenta by signed area interpretation: geometric method (Ricca, 2008)

Let $\mathcal{T} = \bigcup \chi_i$ be a vortex tangle. Then, by considering the associated principal projected graphs, we have

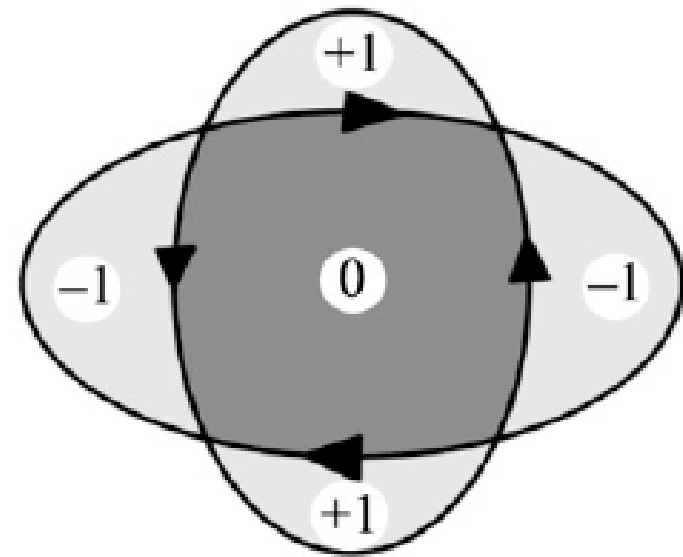
$$\mathbf{P} = (P_{xy}, P_{yz}, P_{zx}) : P_{xy} = \sum_j \gamma_j \mathcal{I}_j \mathcal{A}_{xy}(R_j), \quad P_{yz} = \dots, \quad P_{zx} = \dots ;$$

$$\mathbf{M} = (M_{xy}, M_{yz}, M_{zx}) : M_{xy} = d_z \sum_j \gamma_j \mathcal{I}_j \mathcal{A}_{xy}(R_j), \quad M_{yz} = \dots, \quad M_{zx} = \dots .$$

Examples



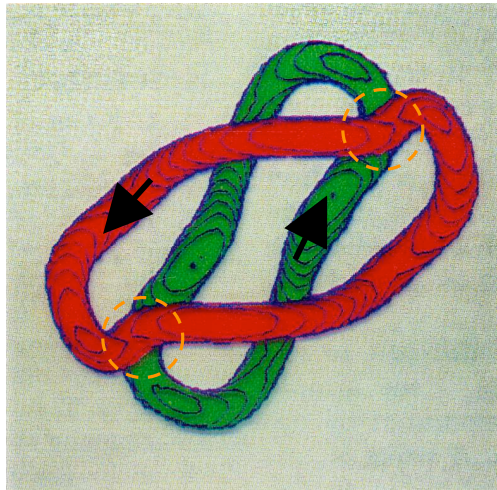
(a)



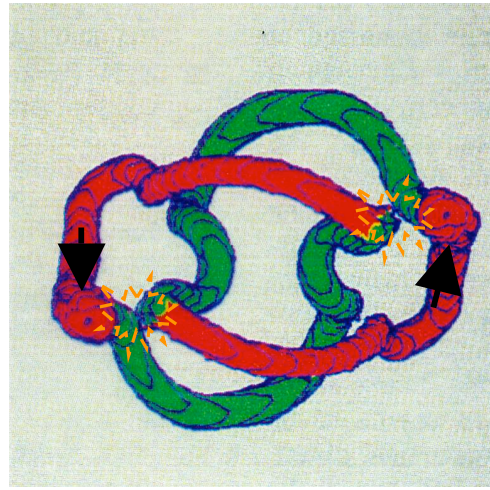
(b)

Vortex analysis by geomertic method

- Following Aref & Zawadzki (Nature, 1991):

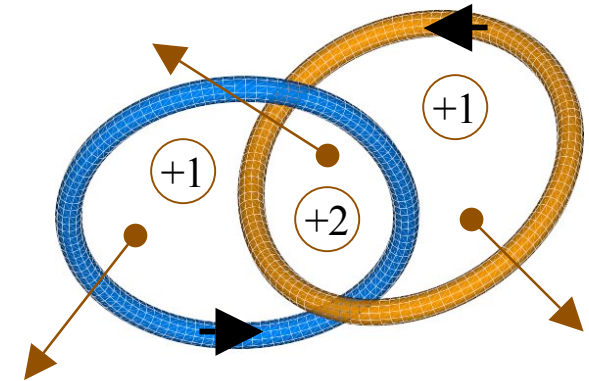


before



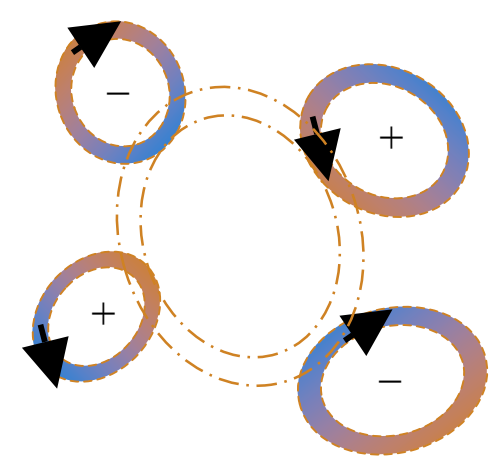
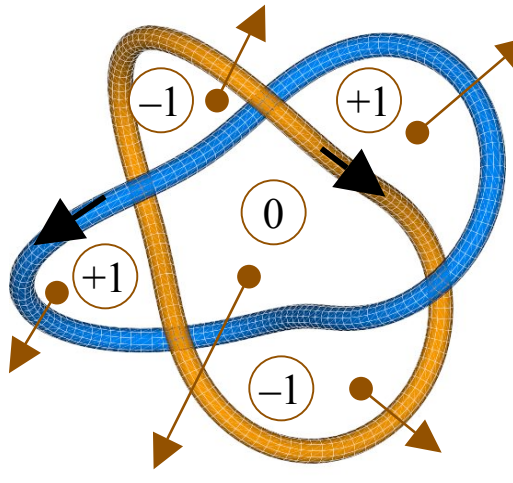
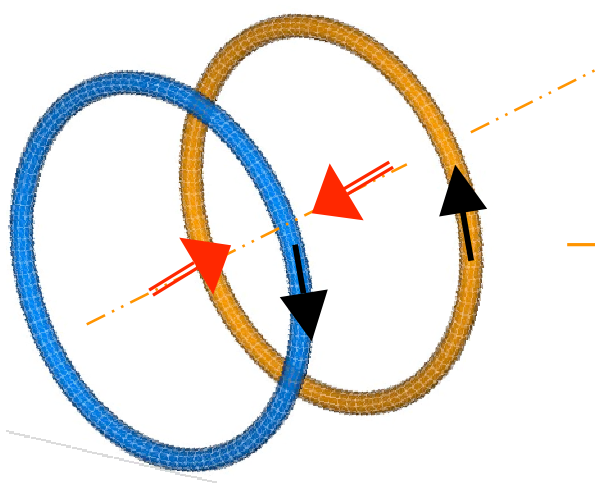
after

...

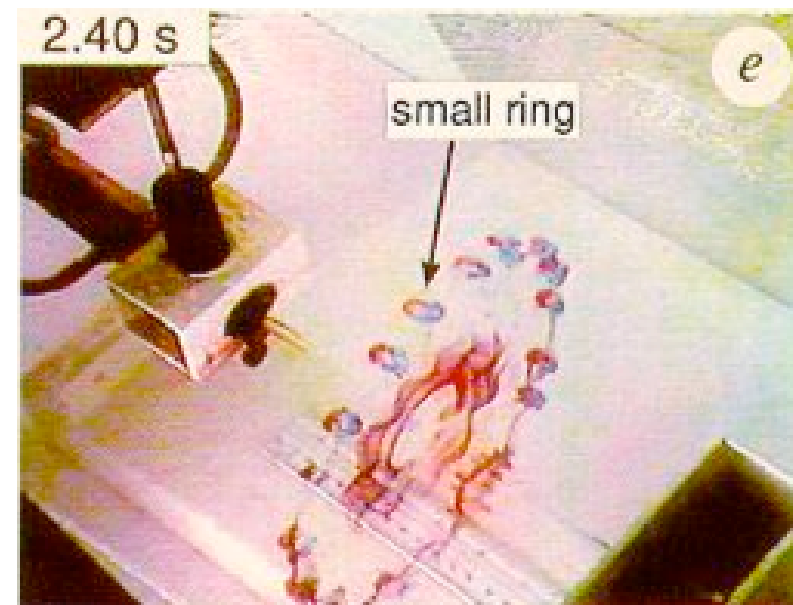
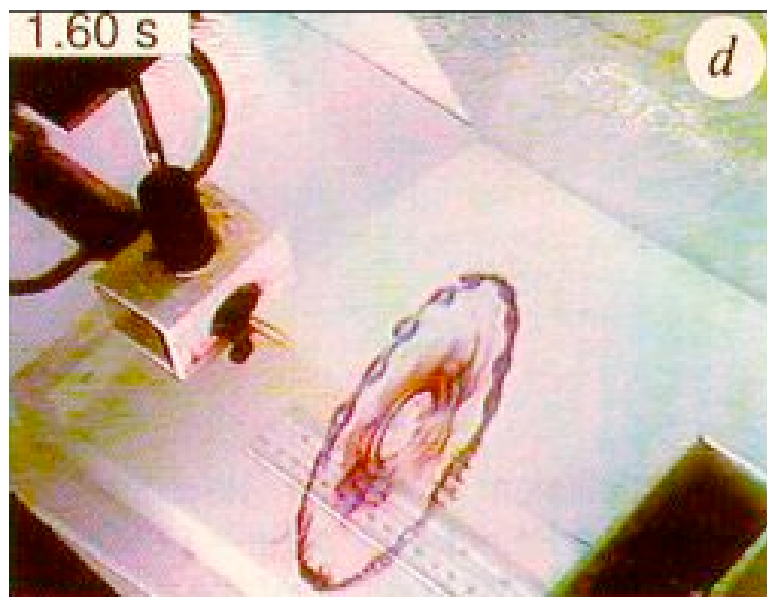
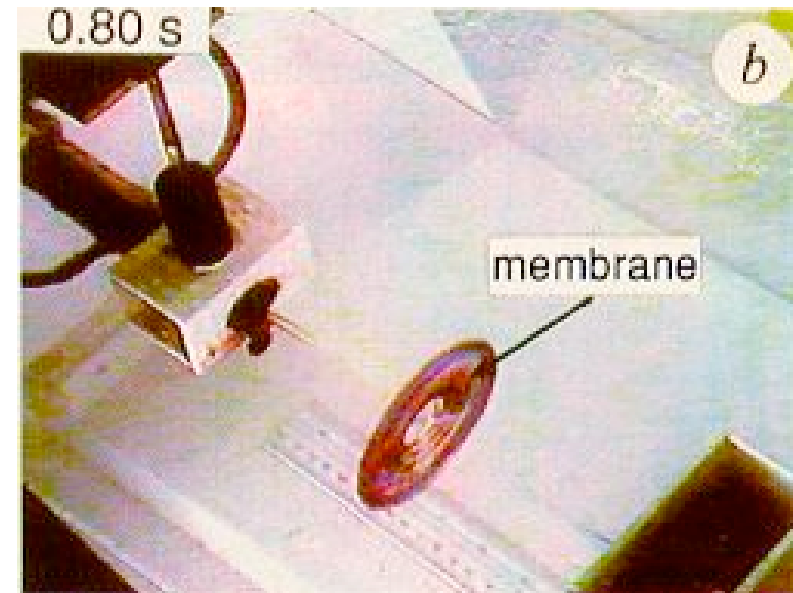
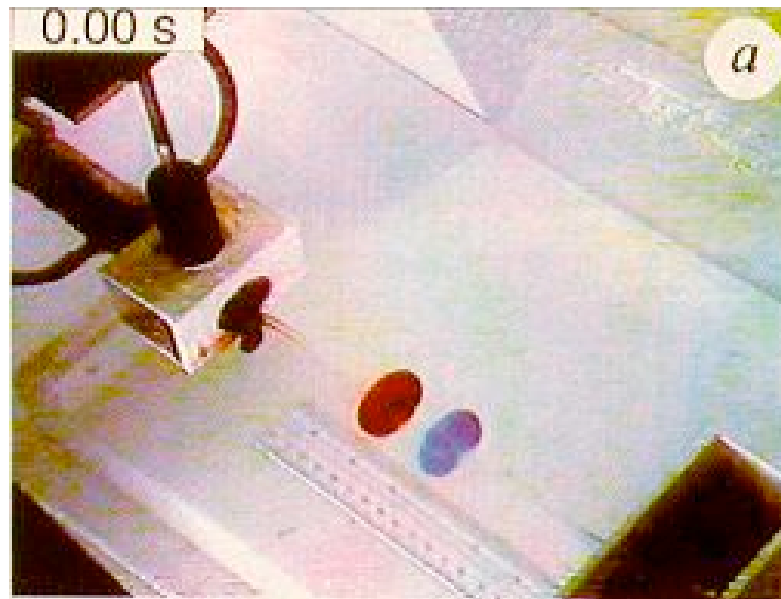


weighted areas

- Zero-area momentum ($P = 0$) interaction:



Head-on collision and breakdown of vortex rings (Lim, Nature, 1992)



Head-on collision of vortex rings (Lim, Nature, 1992): $Re=1071$

Head-on Collision of Coloured Vortex Rings
at $Re=1071$

Note the formation of small rings from the
cross-linking of the wavy vortex filaments
of the larger rings

Head-on collision of vortex rings (Lim, Nature, 1992): $Re=1573$

$Re \approx 1573$

Selected references

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